

# Computer simulation of Poisson's ratio of soft polydisperse discs at zero temperature

J. NAROJCZYK, K. W. WOJCIECHOWSKI\*

Institute of Molecular Physics, Polish Academy of Sciences,  
ul. M. Smoluchowskiego 17, 60-179 Poznań, Poland

A simple algorithm is proposed to study structural and elastic properties of matter in the presence of structural disorder at zero temperature. The algorithm is used to determine the properties of the polydisperse soft disc system. It is shown that Poisson's ratio of the system essentially depends on the size polydispersity parameter; larger polydispersity implies larger Poisson's ratio. In the presence of any size polydispersity, Poisson's ratio increases also when the interactions between the particles tend to the hard potential.

Key words: *Poisson's ratio; polydispersity; elastic constant; inverse-power potential; soft matter*

## 1. Introduction

It is expected that future "intelligent materials" will combine various unusual properties, e.g. unusual electromagnetic properties with unusual elastic properties. Although the elasticity has been the field of human investigation since the ancient times, many problems interesting from both the point of view of basic research and possible applications remain unsolved. One of such problems is the influence of various forms of disorder on the microscopic level on the elastic properties of matter.

In the present note, we concentrate on studies of one of the forms of disorder – the size polydispersity. The question we pose is: How the elastic properties of a system are modified when instead of consisting of identical particles it is formed of particles having some distribution of sizes? Polydisperse systems have recently been intensively studied because of their role in various fields of science and technology [1, 2]. In contrast to the phase diagram and structure, however, their elastic properties remain an open field.

---

\*Corresponding author, e-mail: kww@man.poznan.pl

The subject of our study is a two-dimensional polydisperse model system at zero temperature. The simulations done here were meant to provide additional data to the work described in [3] in which the investigations were carried out at positive temperatures. We concentrate on studies of Poisson's ratio [4] directly describing the deformation of materials under loading/unloading stress. The latter subject is related to increasing interest in so-called auxetic materials, i.e. systems showing negative Poisson's ratio [5–9]. Studies of influence of various mechanisms on Poisson's ratio can help in searching for or manufacturing real materials of this unusual property.

The paper is organized as follows. In Sec. 2, the studied model is described. In section 3, some theoretical background is reminded. In section 4, the chosen method of simulations is sketched. In section 5, the results are presented. Conclusions are drawn in section 6.

## 2. The system studied

The system under study, shown in Fig. 1 and further referred to as the polydisperse soft disc system, consists of soft particles which interact only with their nearest neighbours through the interaction potential of the form

$$u_{ij}(r_{ij}) = \left( \frac{d_i + d_j}{2r_{ij}} \right)^n \quad (1)$$

where  $d_i$  and  $d_j$  are the diameters of the interacting particles. The values of the diameters were generated [10] according to the Gauss distribution function with the fixed size polydispersity parameter  $\delta$ , defined as

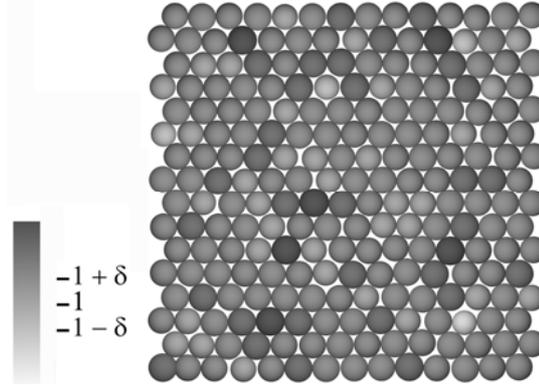
$$\delta = \frac{\left( \langle d^2 \rangle - \langle d \rangle^2 \right)^{1/2}}{\langle d \rangle} \quad (2)$$

It can be seen in Figure 1, that the soft discs form a nearly hexagonal lattice for which the elastic properties are isotropic for small deformations [4]. It is worth to stress that when the exponent  $n$  tends to infinity, the above system tends to the static (i.e., zero temperature), polydisperse, hard disc system, studied in [3] at positive temperatures.

Our aim was to determine the dependence of the elastic constants and Poisson's ratio of a given structure on the amount of disorder introduced into the system; the disorder was quantified by the parameter  $\delta$ . The data which will be shown below were obtained by computer simulation. Although the polydisperse soft disc system is very simple, it includes two main features of various real systems: repulsive forces at high densities and a certain distribution of sizes of the interacting bodies. Hence, one can

expect that it will supply some (qualitative, at least) information on the behaviour of more complex, real systems.

Fig. 1. Typical configuration of the studied system for  $N = 224$  particles;  $\delta$  is the polydispersity parameter. Discs of different sizes are represented by different intensities of greyness



### 3. Theoretical analysis

As the studied system is isotropic in terms of its elastic properties, it can be described by only two elastic constants [4]: the bulk modulus,  $B$ , and the shear modulus,  $\mu$ . The elastic constants  $B$  and  $\mu$  characterize the resistance of material against the changes of the volume and the shape, respectively. Both mentioned values must be positive in the range of the mechanical stability of the system.

The free energy of the deformed two-dimensional (2D) crystal exhibiting a 6-fold symmetry axis can be written as a function of the strain tensor components and the elastic constants [4]:

$$F = -p(\varepsilon_{xx} + \varepsilon_{yy}) + 2\lambda_{\xi\eta\xi\eta}(\varepsilon_{xx} + \varepsilon_{yy})^2 + \lambda_{\xi\xi\eta\eta}[(\varepsilon_{xx} - \varepsilon_{yy})^2 + 4\varepsilon_{xy}^2] \quad (3)$$

The values of the elastic constants and pressure can be obtained by differentiation of the free energy with respect to the strain tensor components, see e.g. [11]:

$$\begin{aligned} \left. \frac{\partial F}{\partial \varepsilon_{xx}} \right|_{\varepsilon=0} &= p \\ \left. \frac{\partial^2 F}{\partial \varepsilon_{xx}^2} \right|_{\varepsilon=0} &= 4\lambda_{\xi\eta\xi\eta} + 2\lambda_{\xi\xi\eta\eta} \\ \left. \frac{\partial^2 F}{\partial \varepsilon_{xx} \partial \varepsilon_{yy}} \right|_{\varepsilon=0} &= 4\lambda_{\xi\eta\xi\eta} - 2\lambda_{\xi\xi\eta\eta} \end{aligned} \quad (4)$$

The bulk modulus,  $B$ , and the shear modulus,  $\mu$ , can be easily related [12] to the quantities used in Eq. (3):

$$B = 4\lambda_{\xi\eta\xi\eta}, \quad \mu = 2\lambda_{\xi\xi\eta\eta} - p \quad (5)$$

Poisson's ratio  $\nu$  for an isotropic 2D system can be expressed by the elastic moduli as follows [12]:

$$\nu = \frac{B - \mu}{B + \mu} \quad (6)$$

In the case of the discs of equal diameters ( $\delta = 0$ ), one can obtain analytical formulae for the elastic constants:

$$p_0 = n\sqrt{3}a^{-n-2}, \quad B_0 = \left(\frac{n}{2} + 1\right)p_0, \quad \mu_0 = \left(\frac{n}{4} - \frac{1}{2}\right)p_0, \quad \nu_0 = \frac{n+6}{3n+2} \quad (7)$$

where  $n$  is the power of the potential and  $a$  is the distance between the centres of particles.

#### 4. The method of simulation

The elastic properties with respect to the degree of polydispersity [3] were studied. For each of more than 1000 different structures the reference state of minimal energy was found. Then the deformation and the "measurement" took place.

Simulations were based on the algorithm searching for a configuration of minimum energy for a given shape of the system. The model was enclosed in periodic boundary conditions. Its total energy was a sum of the interactions of the form (1) between the nearest-neighbouring particles (i.e., those whose Dirichlet polygons have a common side). In search for the minimum energy states, the program conducted a series of translational moves of particles; only those decreasing the energy of the system were accepted. The particles were moved in the direction of the total force acting on them by an arbitrary vector  $d\mathbf{r}_i$ . After acceptance/rejection of each move, the vector was modified with respect to the change of the total energy of the system: if the move led to increase/decrease of the energy  $d\mathbf{r}_i$  was decreased/increased, respectively. The particles were moved one-by-one in the order they were placed in the lattice. Enumeration of the particles was arbitrary and was checked to have no influence on the application of the algorithm and its results. To avoid trapping the system into a local minimum, some "shaking" (random changes of the particle positions) was introduced. Each particle was randomly displaced by a vector  $d\mathbf{r}_{\text{rand}}$  of a random length and orientation, generated within a defined range of values. The program moved all the particles of the system as long as the system energy was no longer changing within the given accuracy.

After calculating the energy of the system at equilibrium (equal to its free energy at zero temperature) the shape of the system was deformed by strains for which the minimum energy configurations of particles were searched for. Then a numerical differentiation of the energy with respect to the deformations was used to determine the

values of elastic constants  $B$  and  $\mu$  from which Poisson's ratio of the system was obtained. The precise description of the algorithm will be given elsewhere.

### 5. Results

The figures presented below summarize the results of simulation. It can be seen that for large values of the exponent  $n$ , the elastic properties of the system strongly depend on its polydispersity.

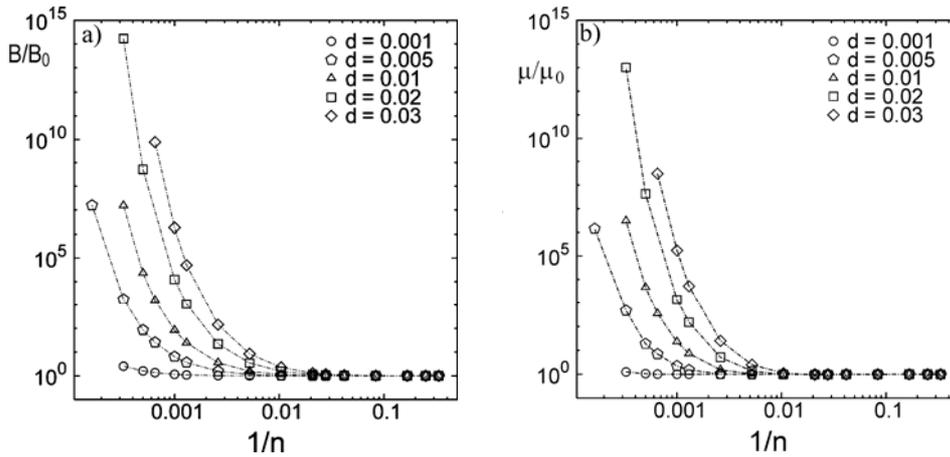


Fig. 2. The dependence of: a) the bulk modulus  $B$ , b) the shear modulus  $\mu$ , on the inverse of the exponent  $n$ . The dotted lines are presented to guide the eye

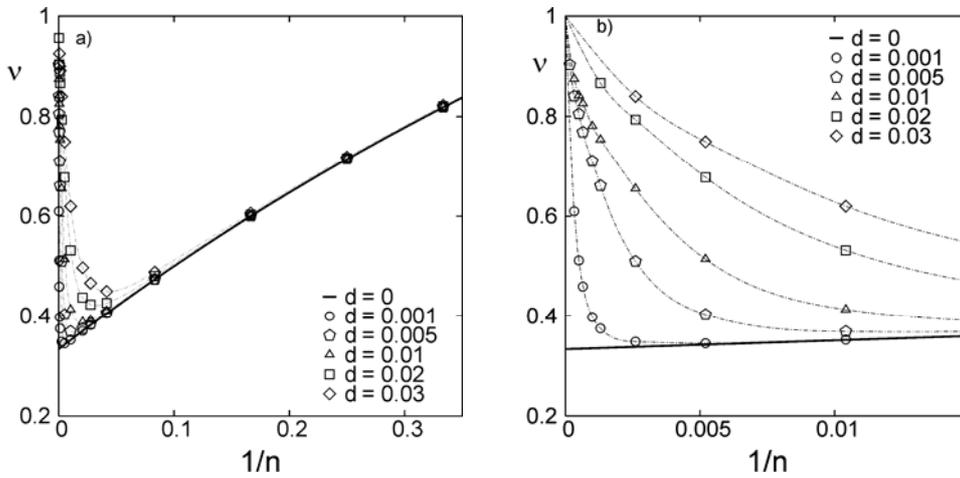


Fig. 3. The value of Poisson's ratio against the inverse of the exponent  $n$  for two-dimensional soft discs with different values of the polydispersity parameter: a)  $n \geq 3$ , b)  $n \geq 96$ . The thick solid line shows the exact result for  $\delta = 0$ , see Eq. (7). The dotted lines are presented to guide the eye

In particular, it can be seen in Fig. 2 that when  $n \rightarrow \infty$ , both  $B$  and  $\mu$  grow rapidly with respect to the values obtained for the systems composed of identical particles. By increasing the polydispersity of the system, its Poisson's ratio grows with respect to the periodic system of discs with identical diameters (Fig. 3). When  $n \rightarrow \infty$ , then for any non-zero polydispersity, Poisson's ratio tends to its maximum value possible for a two-dimensional isotropic system equal 1. The elastic properties of the polydisperse soft disc system at zero temperature are in very good agreement with the Monte Carlo simulations of this system at positive temperatures, when the temperature tends to zero [13].

## 6. Conclusions

A simple and efficient algorithm to study static structures of polydisperse soft particles has been proposed. It was shown that the static disorder studied (the lack of periodicity of the particle positions caused by unequal sizes of the particles) has an essential influence on the value of Poisson's ratio. Typically, for positive values of the polydispersity parameter, the bulk modulus  $B$ , the shear modulus  $\mu$ , and Poisson's ratio  $\nu$ , increase with increasing  $n$  when potentials become strongly repulsive, i.e. when the exponent  $n$  is large, which corresponds to the hard particle limit. These quantities increase also with increasing polydispersity parameter  $\delta$ , in the same limit of large values of  $n$ .

The results obtained in this work show an excellent agreement with recent results obtained by the Monte Carlo simulations of the soft polydisperse discs at the low temperature limit [3], which, in turn, agree very well with the results obtained for hard polydisperse discs [13] in the limit of high temperatures and large  $n$ . This confirms the hypothesis that the elastic properties of hard-body systems can be obtained by studying soft-body systems and taking the  $n \rightarrow \infty$  limit. We plan to apply the method to three-dimensional systems of isotropic and anisotropic particles recently studied at positive temperatures. We expect that studies of very simple models, as the one described in the present paper, will help in understanding and predicting behaviour of real systems. It is interesting to check to what extent such simple models approximate the properties of some real systems like those described in [14].

### Acknowledgements

This work was supported by the Polish Committee for Scientific Research (grant No. 4T11F01023) and by the Center of Excellence for Magnetic and Molecular Materials for Future Electronics within the European Commission contract No. G5MA-CT-2002-04049. The simulations were performed at the Poznań Supercomputing and Networking Centre.

### References

- [1] SOLLICH P.J., *J. Phys.: Condens. Matter*, 14 (2002), R79.
- [2] XU H., BAUS M.J., *J. Chem. Phys.*, 118 (2003), 5045.

- [3] TRETIKOV K.V., WOJCIECHOWSKI K.W., Phys. Status Sol. B, 242 (2005), 730.
- [4] LANDAU L.D., LIFSHITS E.M., *Theory of Elasticity*, Pergamon Press, London, 1986.
- [5] LAKES R.S., Science, 235 (1987), 1038; *ibid* 238 (1987), 551 .
- [6] WOJCIECHOWSKI K.W., Phys. Lett., A 137 (1989), 60.
- [7] EVANS K.E., ALDERSON A., Adv. Mater., 12 (2000), 617.
- [8] WOJCIECHOWSKI K.W., *Monte Carlo simulations of model particles forming phases of negative Poisson's ratio*, [in:] *Properties and Applications of Nanocrystalline Alloys from Amorphous Precursors*, B.Idzikowski, P. Svec, M. Miglierini (Eds.), Kluwer, Dordrecht, 2005, 241–252.
- [9] WOJCIECHOWSKI K.W., ALDERSON K.L., BRAŃKA A.C., ALDERSON K.L., Phys. Status Sol., B 242 (2005), 497; <http://www.ifmpan.poznan.pl/zp10/auxet2/main.html>.
- [10] PRESS W.H., TEUKOLSKY S.A., VETTERLING T.W., FLANNERY B.P., *Numerical Recipes in FORTRAN*, 2nd Ed., Cambridge University Press, Cambridge, 1992.
- [11] WOJCIECHOWSKI K.W., TRETIKOV K.V., KOWALIK M., Phys. Rev. E, 67 (2003), 0361211.
- [12] WOJCIECHOWSKI K.W., J. Phys. A: Math. Gen., 36 (2003), 11765.
- [13] TRETIKOV K. V., WOJCIECHOWSKI K. W., Phys. Rev. E, unpublished.
- [14] VAN BLAADEREN A., MRS Bulletin, 29, (2004), 85.

*Received 15 September 2005*

*Revised 30 November 2005*